

OBSERVATIONS OF TYPE Ia SUPERNOVA 2014J WITH FLITECAM ON SOFIA

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ABSTRACT

We present medium-resolution near-infrared (NIR) spectra, covering 1.1–3.4 μm , of the normal Type Ia supernova (SN Ia) SN 2014J in M82 obtained with the FLITECAM instrument on board Stratospheric Observatory for Infrared Astronomy (SOFIA) between 17 and 26 days after maximum B light. Our 2.8–3.4 μm spectra may be the first $\sim 3 \mu\text{m}$ spectra of an SN Ia ever published. The spectra spanning the 1.5–2.7 μm range are characterized by a strong emission feature at $\sim 1.77 \mu\text{m}$ with a FWHM of $\sim 11,000\text{--}13,000 \text{ km s}^{-1}$. We compare the observed FLITECAM spectra to the recent non-LTE delayed detonation models of Dessart et al. and find that the models agree with the spectra remarkably well in the 1.5–2.7 μm wavelength range. Based on this comparison we identify the $\sim 1.77 \mu\text{m}$ emission peak as a blend of permitted lines of Co II. Other features seen in the 2.0–2.5 μm spectra are also identified as emission from permitted transitions of Co II. However, the models are not as successful at reproducing the spectra in the 1.1–1.4 μm range or between 2.8 and 3.4 μm . These observations demonstrate the promise of SOFIA, which allows access to wavelength regions inaccessible from the ground, and serve to draw attention to the usefulness of the regions between the standard ground-based NIR passbands for constraining SN models.

Key words: infrared: stars – supernovae: general – supernovae: individual (SN 2014J)

1. INTRODUCTION

It is widely accepted that Type Ia supernovae (SNe Ia) originate from the runaway thermonuclear burning of a carbon-oxygen white dwarf that approaches or exceeds the Chandrasekhar mass limit. However, the evolutionary scenario and the mechanism by which the white dwarf gains mass, e.g., accretion from a binary companion in the so-called single-degenerate model (Hoyle & Fowler 1960; Whelan & Iben 1973) or merger with another white dwarf in a double-degenerate model (Iben & Tutukov 1984; Webbink 1984) that trigger the ignition and give rise to the SN, as well as the details of the explosion process itself, remain uncertain. Given the importance of SN Ia to distance determinations and cosmology (see, e.g., Goobar & Leibundgut 2011), it is crucial to understand the physics behind these explosions. One of the best ways to do this is to carry out detailed comparisons between high quality observational spectra of bright SNe Ia with physical models. Near-infrared (NIR) spectroscopic observations of SNe Ia are particularly important in this regard as they are not significantly affected by the presence of foreground or circumstellar dust that can seriously impact and complicate the interpretation of ultraviolet and optical spectra. In addition, the line blanketing opacity that dominates the UV

and optical regions of SNe Ia spectra is lower in the NIR, making the interpretation and attribution of features easier. In order to reduce uncertainties in the model parameters and improve the accuracy of the models themselves, continuous spectra of bright nearby SNe Ia are necessary against which the models can be compared. This includes the regions between the standard ground-based NIR pass-bands that are inaccessible from most observatories due to strong atmospheric absorption. In addition, continuous spectra spanning these telluric absorption regions are particularly important for deriving k -corrections, which are needed to properly compare photometric observations of SNe Ia at high redshift with those of lower redshift counterparts (e.g., Hsiao et al. 2007, and references therein). These corrections account for the differences in the rest-frame wavelength regions sampled by fixed wavelength filters as cosmic expansion shifts and stretches the intrinsic spectrum of an object. Gaps in the wavelength coverage of low-redshift template objects, such as those resulting from telluric absorption, increase the uncertainties in the derived corrections, making comparisons between high-redshift and low-redshift samples difficult (see e.g., Hsiao et al. 2013).

SN 2014J in M82 was discovered on 2014 January 21 by Fossey et al. (2014). Zheng et al. (2014) estimated the time of

first light to be January 14.75 UT, about 7 days prior to discovery, while Marion et al. (2014) derived the time of maximum B brightness to be February 01.74 UT. One of the nearest and brightest SNe Ia in modern history, SN 2014J has been studied across a large range of wavelengths since its discovery, allowing the determination of a number of its physical parameters. Ultraviolet, optical, NIR, and mid-infrared observations have revealed that SN 2014J was a spectroscopically normal SN Ia, nearly identical to SN 2011fe, although with substantially higher reddening ($A_V \sim 2.0$; Amanullah et al. 2014; Foley et al. 2014; Goobar et al. 2014; Marion et al. 2014; Zheng et al. 2014) and relatively higher velocity ($v \geq 11,000 \text{ km s}^{-1}$) absorption features in its optical and NIR spectra (Goobar et al. 2014; Marion et al. 2014). The analysis of the optical and NIR spectra presented by Marion et al. (2014) reveals radial stratification of elements produced during the explosion, indicative of relatively little mixing and consistent with predictions from delayed detonation models. Telesco et al. (2014) were able to reproduce successfully the 8–13 μm spectra of SN 2014J, which are dominated by strong lines of [Co II] and [Co III], with a model from Höflich et al. (2002) for a spherical delayed detonation explosion on a Chandrasekhar-mass white dwarf. Based on fits to the observed gamma-ray line fluxes, Churazov et al. (2014) and Diehl et al. (2015) estimated that $\sim 0.5\text{--}0.6 M_\odot$ of ^{56}Ni was produced in the explosion, in agreement with the model used by Telesco et al. (2014).

Because SN 2014J occurred in a nearby galaxy ($D \sim 3.4 \text{ Mpc}$; Dalcanton et al. 2009) and therefore was relatively bright, it constitutes a nearly ideal object on which to test SN Ia spectral models. Fortuitously, it also provided an opportunity to test the FLITECAM instrument (McLean et al. 2006) on the NASA Stratospheric Observatory for Infrared Astronomy (SOFIA), which can be used to fill in some of the wavelength gaps in NIR spectra of SNe Ia obtained from ground-based facilities. NIR observations obtained with SOFIA are much less susceptible to the deleterious effects of atmospheric water vapor, whose strong absorptions serve to define the standard ground-based J ($\sim 1.1\text{--}1.35 \mu\text{m}$), H ($\sim 1.5\text{--}1.75 \mu\text{m}$), and K ($\sim 2.0\text{--}2.35 \mu\text{m}$) bands. At altitudes $>39,000$ feet ($\sim 12 \text{ km}$), typical of SOFIA flights, the precipitable water vapor overburden is generally between a few to tens of microns whereas ground-based precipitable water vapor overburdens are typically a factor of 1000 larger, on the order of a few millimeters, even at the highest and driest sites (e.g., Giovanelli et al. 2001; Radford et al. 2008). A description of SOFIA and its instrument suite can be found in the papers by Young et al. (2012), Gehrz et al. (2009), and Becklin et al. (2007). In this paper we present the first spectra obtained by FLITECAM during its commissioning flights on board SOFIA. The NIR spectra of SN 2014J obtained with FLITECAM/SOFIA and discussed here comprise a complement to the data sets presented by Marion et al. (2014) and Friesen et al. (2014).

2. OBSERVATIONS AND DATA REDUCTION

SN 2014J was observed with FLITECAM on four separate SOFIA flights (flight nos. 146, 147, 148, and 149) on 2014 February 19, 21, 25, and 27 UT, approximately 17–25 days after the time of maximum B brightness, or approximately 36–44 days after first light (Zheng et al. 2014). The flights originated in Palmdale, CA and the primary goal was the commissioning of the FLITECAM instrument on the SOFIA

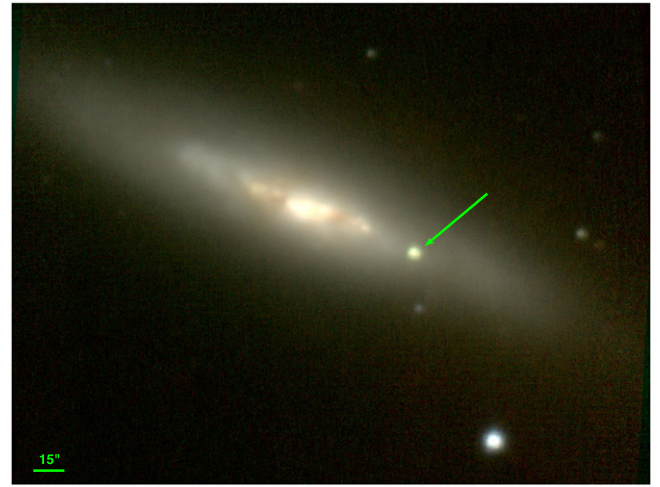


Figure 1. Three color image of the SN 2014J field obtained with FLITECAM. Blue is J , green is H , and red is K . The arrow identifies SN 2014J, which is readily apparent as the point source SW of the nucleus and in the plane of the disk of M82. North is up and east is left.

telescope. FLITECAM is a NIR imager and grism spectrograph covering the 1–5 micron range (McLean et al. 2006; Smith & McLean 2008; Logsdon et al. 2014) and both images and spectra of the SN were obtained. FLITECAM’s 1024×1024 InSb ALADDIN III array covers an $\sim 8'$ diameter field of view with a plate scale of $0''.475 \text{ pix}^{-1}$. Dithered images were obtained in the FLITECAM J , H , and K filters.¹⁶ FLITECAM was co-mounted with the HIPO instrument during these observations, a configuration that precluded observations at wavelengths longwards of $\sim 4 \mu\text{m}$, and reduced the sensitivity longwards of $\sim 2 \mu\text{m}$, due to high background levels resulting from the warm dichroic and transfer optics. The observations were acquired at four different altitudes between 38,000 and 43,000 feet and a range of air masses. The FWHM of the images and spectra on these flights was $\sim 3''$. The imaging observations were reduced in a manner typical for NIR images. Flat fields and background sky frames were constructed from dithered observations obtained $6'$ away from the target. Individual frames were flat-fielded, sky-subtracted, and then shifted and median combined to produce a final image in each filter. The frames in each filter were then aligned and combined (see Figure 1). Photometric calibrations of the images were determined from measurements of the star BD+70 587, which is present on the FLITECAM images of SN 2014J and for which 2MASS magnitudes are available. The resulting SN photometry is presented in Table 1. The agreement with the values given by Amanullah et al. (2014), Marion et al. (2014), and Foley et al. (2014) for times near those of our SOFIA observations is reasonably good.¹⁷

For spectroscopic observations, FLITECAM contains three grisms, each of which can be used in three orders and with various order sorting filters, to provide wavelength coverage between ~ 1.1 and $\sim 5.5 \mu\text{m}$ at medium resolution. Two slits,

¹⁶ The FLITECAM filter passbands are available at <http://www.sofia.usra.edu/Science/ObserversHandbook/FLITECAM.html>.

¹⁷ We note that because the spectrum of SN 2014J is unlike that for any standard star, the shifts in the effective wavelengths of the filters and therefore the color corrections could be substantial, particularly in the J and H bands. However, in order to facilitate comparisons with magnitudes of the SN given in the literature, we have not made any such corrections to the values given in Table 1.

Table 1
FLITECAM/SOFIA Photometric Observations of SN 2014J

MJD	UT Date	Phase (days)	Filter	Mag
56707.3	2014 Feb 19 07:11	17.5	J	10.97 ± 0.06
56707.3	2014 Feb 19 07:19	17.5	H	9.64 ± 0.03
56707.3	2014 Feb 19 07:29	17.5	K	9.54 ± 0.04
56709.3	2014 Feb 21 07:22	19.5	J	10.93 ± 0.05
56709.3	2014 Feb 21 07:32	19.5	H	9.56 ± 0.04
56709.3	2014 Feb 21 07:43	19.5	K	9.49 ± 0.04

Note. Phase is the number of days past the time of maximum *B* light (MJD 56689.8; Marion et al. 2014).

each 60'' long, are available and provide spectral resolving powers R of ~ 1300 and ~ 1700 , respectively. Spectra are usually obtained by nodding the target between two different positions along the slit, referred to as A and B respectively. Due to the fast rotation of the field and the constraints of pointing the telescope, the spectra of SN 2014J were acquired via several sequences of AB observations. The A and B positions were separated by 30''. The wider, low-resolution slit was used for the observations, which resulted in $R \sim 1300$ for the SN spectra. The log of the spectroscopic observations is given in Table 2. The data were reduced using Redux, the SOFIA facility pipeline (Clarke et al. 2014) incorporating the fspextool software package, which is modification of the Spextool package (Cushing et al. 2004) developed for the SpeX NIR spectrograph at the NASA Infrared Telescope Facility (Rayner et al. 2003). fspextool performs non-linearity correction, AB pair subtraction, source profile construction, extraction and background aperture definition, optimal extraction, and wavelength calibration for FLITECAM grism data. Wavelength calibration was performed using the OH emission lines in the spectra of the background sky.

In order to account for the differences in the detector response at the two positions (A and B) along the slit, we did not combine the A and B spectra until after the telluric correction and flux calibration steps. The latter steps were carried out using the xtellcor_general package (Vacca et al. 2003), and for these purposes we observed an A0V star during each flight. The spectra of the A0V star were reduced in exactly the same way as for the SN. Again, we kept the A and B beam spectra separate. We flux calibrated the A beam spectra of the SN with the A beam spectra of the standard and did the same for the B beam spectra. To remove residual telluric features resulting from the differences in observing altitude and air mass between the standard star and the SN, we computed telluric absorption models appropriate for the observing conditions, using the ATRAN code (Lord 1992), smoothed and binned them to the FLITECAM resolution and sampling, and generated correction curves as a function of wavelength from the ratios of the telluric models for the standard to those for the SN. These corrections were then applied to the spectra of the SN. The telluric-corrected and flux-calibrated A and B beam spectra of the SN were then combined. The spectra from each grism setting for a given date were then merged to produce a final spectrum spanning the full wavelength range available. We then scaled the spectra so that synthetic photometry yielded values that matched the imaging photometry values provided by our own images and those of Foley

et al. (2014), Amanullah et al. (2014), and Marion et al. (2014) interpolated to the dates of our observations.

3. RESULTS

The extracted FLITECAM spectra of SN 2014J are shown in Figure 2. Wavelength regions that are inaccessible, or difficult to observe, from the ground due to strong telluric absorption are marked. In Figure 3 we present a comparison between the spectra from February 19 and February 27 over the 1.5–2.7 μm wavelength range. The signal-to-noise ratio (S/N) for these spectra ranges from ~ 10 to more than 60 per pixel at the highest flux levels; over most of the wavelength range, the S/N is > 20 per pixel.

The most striking aspects of the spectra shown in Figures 2 and 3 are the strong emission features at 1.55–1.65 μm and $\sim 1.77 \mu\text{m}$. The former emission feature can clearly be seen in the later-time ground-based spectra of SN 2014J presented by Marion et al. (2014). Although the latter feature is present in ground-based spectra of SN 2014J (Marion et al. 2014), and can also be seen in some of the spectra of other SN Ia presented by Gall et al. (2012), Hsiao et al. (2013), and Marion et al. (2003, 2009), the FLITECAM spectra are the only data that span the entire feature profile; the ground-based spectra cover only a portion of this wavelength region due to the poor and variable atmospheric transmission between ~ 1.8 and $\sim 2.0 \mu\text{m}$. The shapes and peak wavelengths of both features vary substantially between the two sets of observations, which were obtained only ~ 8 days apart. A Gaussian fit to the 1.77 μm feature in the February 19 spectrum yields a centroid wavelength of 1.768 μm and a FWHM of $\sim 0.080 \mu\text{m}$ ($\sim 13,500 \text{ km s}^{-1}$); on February 27, the feature is centered at 1.784 μm with a width of $\sim 0.062 \mu\text{m}$ ($\sim 10,500 \text{ km s}^{-1}$). Uncertainties are on the order of 0.003 μm ($\sim 500 \text{ km s}^{-1}$). If the observed wavelength shift were to reflect a change in velocity of a specific emission feature, this would correspond to a velocity redshift of about 2700 km s^{-1} , with a decrease in the velocity width of a similar amount. Although the ground-based data presented by Marion et al. (2014) confirm the shift seen in the FLITECAM spectra, the full wavelength coverage provided by FLITECAM in the 1.8–2.0 μm region allows the change in the profile to be identified easily. A shift in the feature wavelength is not readily apparent in the NIR spectra of other SNe Ia shown by Marion et al. (2009) but can be discerned in the spectra of the normal SN Ia SN 2011fe presented by Hsiao et al. (2013).

Based on the model of Wheeler et al. (1998), Marion et al. (2003, 2009) identify the emission in the 1.7 μm region as arising from a blend of Fe group lines along with absorption from Co II. However, the strength and symmetric Gaussian shape of the $\sim 1.77 \mu\text{m}$ feature argue against this interpretation and in favor of (primarily) a single emission feature. Dessart et al. (2014) identify this feature, as well as other strong features seen in the *K* band spectra of SNe Ia, as emission from permitted lines of Co II. In fact, the strengthening of the Co lines is responsible for the well-known secondary peak in the broadband NIR photometric light curves of SNe Ia in their model (see also Jack et al. 2012). Their DDC10 A4D1 (hereafter DDC10) model spectra exhibit a strong Co II emission feature at $\sim 1.75 \mu\text{m}$, arising from a blend of lines, primarily those at 1.7463, 1.7770, and 1.8069 μm , starting at about 25 days after explosion and shifting redward with time. The redshift is the result of the variation in line strengths of the

Table 2
Log of FLITECAM/SOFIA Spectroscopic Observations of SN 2014J

MJD	UT Date	$t-t_0$ (days)	$t-t_{B\max}$ (days)	Grism	Wavelengths (μm)	t_{exp} (s)	N_{exp}	Altitude (feet)	ZA
56707.4	2014 Feb 19 08:52	35.6	17.6	H_C	1.500–1.718	300	4	41,000	41
56707.4	2014 Feb 19 10:26	35.7	17.6	Hw_B	1.675–2.053	150	8	41,000	37
56707.5	2014 Feb 19 10:50	35.7	17.7	Kw_C	1.910–2.276	150	8	41,000	41
56707.5	2014 Feb 19 11:28	35.7	17.7	Kl_A	2.270–2.722	180	2	41,000	43
56707.5	2014 Feb 19 11:35	35.7	17.7	Kl_A	2.270–2.722	120	2	41,000	43
56709.4	2014 Feb 21 09:35	37.7	19.6	J_B	1.141–1.385	180	4	43,000	40
56709.4	2014 Feb 21 09:53	37.7	19.6	Kl_A	2.270–2.722	120	6	43,000	40
56713.4	2014 Feb 25 10:17	41.7	23.6	LM_C	2.779–3.399	150	4	40,000	35
56713.4	2014 Feb 25 10:31	41.7	23.6	LM_C	2.779–3.399	200	12	40,000	35
56715.3	2014 Feb 27 06:20	43.5	25.5	J_B	1.141–1.385	180	4	38,000	34
56715.3	2014 Feb 27 06:37	43.5	25.5	H_C	1.500–1.718	180	4	38,000	33
56715.3	2014 Feb 27 06:53	43.5	25.5	Hw_B	1.675–2.053	180	4	38,000	33
56715.3	2014 Feb 27 07:10	43.6	25.5	Kl_A	2.270–2.722	180	4	38,000	33
56715.3	2014 Feb 27 07:27	43.6	25.5	Kw_C	1.910–2.276	180	4	38,000	33

Note. UT date and time given are at the beginning of an exposure set. The exposure time t_{exp} is the time for a single exposure, and N_{exp} is the number of exposures obtained. $t-t_0$ is the number of days past t_0 , the first light date estimated by Zheng et al. (2014); $t-t_{B\max}$ is the number of days past maximum B light estimated by Marion et al. (2014). ZA is the average zenith angle of the target during the observations.

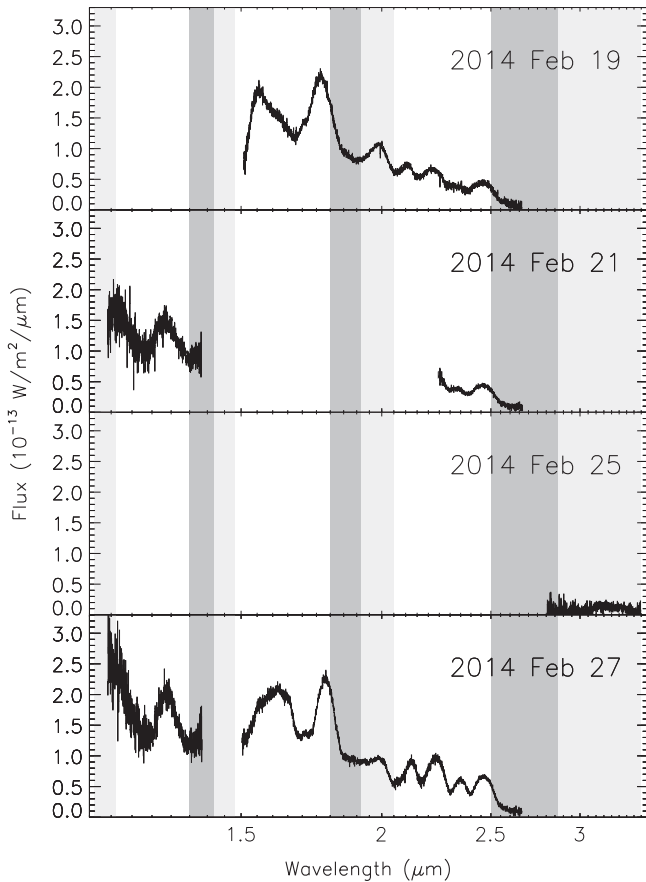


Figure 2. FLITECAM flux-calibrated spectra of SN 2014J on the four different observing dates. Regions of strong telluric absorption from the ground (transmission <20%) are shown in dark gray, while regions of moderate telluric absorption from the ground (transmission <80%) are shown in light gray (cf., for example, Rayner et al. 2009).

various components. Similarly, Gall et al. (2012) identify both the 1.55 μm and the 1.77 μm feature seen in the NIR spectra of the normal SN Ia SN 2005cf and three other SNe Ia at 9–12 days past maximum B -band emission, as permitted lines of Co II, based on comparisons with the W7 model of Nomoto et al. (1984).

In Figure 3 we compare our combined 1.5–2.7 μm spectra of SN 2014J at 18 and 26 days past B maximum (36 and 44 days past first light) with the DDC10 model spectra¹⁸, computed for dates close to those of our FLITECAM observations. The DDC10 A4D1 model is a set of one-dimensional (i.e., spherically symmetric), time-dependent, non-LTE calculations for a delayed detonation explosion of a Chandrasekhar-mass white dwarf at a distance of 10 pc. The model includes non-thermal effects, non-local energy deposition, “huge” model atoms for iron group elements, and forbidden lines of both intermediate mass elements and iron group elements; descriptions of the model can be found in Dessart et al. (2014), Blondin et al. (2013), and references therein. For these comparisons, we have assumed that the time of the SN explosion is given by the time of first light (2014 January 14.75 UT) estimated by Zheng et al. (2014). For these models the difference between the two times should be less than a day (S. Blondin, private communication). We shifted the model spectra to account for the redshift of M82 ($v = 203 \text{ km s}^{-1}$). We also reddened the model spectra to account for both foreground reddening due to the Milky Way ($A_V = 0.06$) and circumstellar reddening and scattering in the vicinity of the SN itself, using the extinction value and the power law reddening/scattering curve advocated by Amanullah et al. (2014; $A_V = 1.77$ and $p = -2.1$; see also Foley et al. 2014). We then scaled the model spectra to match the overall flux levels of the observed data; a single scale factor for each model spectrum was derived from a

¹⁸ Kindly provided to us by L. Dessart and S. Blondin.

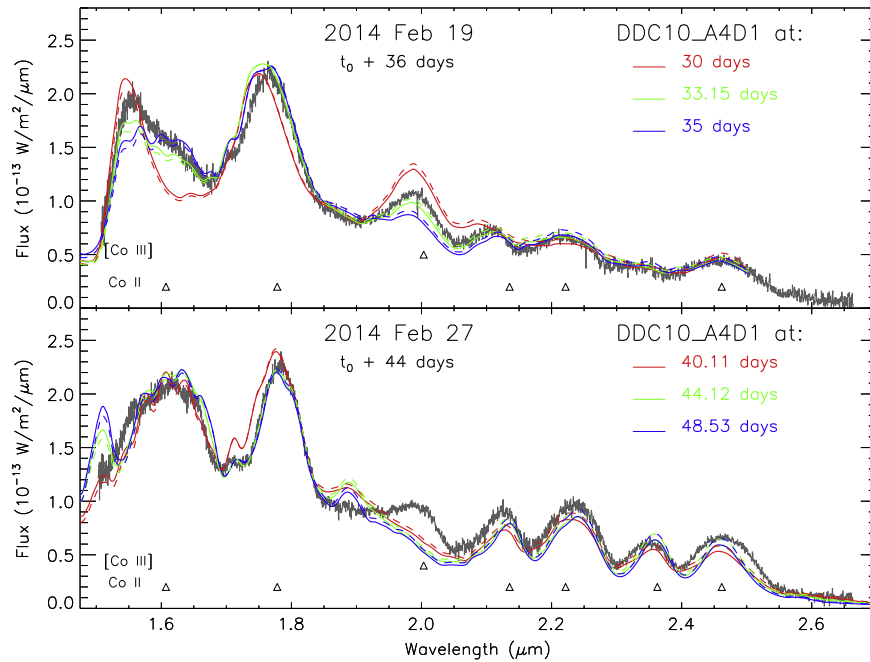


Figure 3. Comparison between the observed $H + K$ -band spectrum of SN 2014J obtained with FLITECAM (in gray) on 2014 February 19 UT (top) and 2014 February 27 UT (bottom) and the CMFGEN model DDC10 A4D1 of Dessart et al. (2014) for three different times past explosion. The model spectra have been scaled to match the flux levels of the data. Dashed lines are model spectra that have been reddened due to the effects of dust in both the Milky Way and M82. Triangles denote the locations of lines of Co II and [Co III] predicted to be strong in the model.

least-squares fit to the observed spectra. The strongest emission lines of Co II and [Co III], as predicted by the model (Blondin et al. 2015), are identified near the bottom of each plot. (The wavelengths of NIR Co II lines can be found in the tables given by Marion et al. 2009 and Gall et al. 2012.)

As can be seen, the model spectra match the observed spectra extremely well in this wavelength regime, although the time past explosion for the best-fitting models for the two observed spectra is somewhat discrepant. The model spectrum that best matches the FLITECAM spectra obtained on February 19 (~ 36 days after first light) corresponds to 33.15 days past explosion, whereas the model spectrum corresponding to 40.11 days past explosion provides the formal best fit to the FLITECAM spectra obtained on February 27 (~ 44 days past first light). In the latter case, however, the model spectra for 44.12 and 48.53 days past explosion fit the observed data in the region around the $1.77 \mu\text{m}$ feature considerably better.

Since the model spectra are computed for a distance of 10 pc, the scale factors needed to match the model spectra to the observed flux levels yield estimates of the distance to SN2014J and therefore provide additional constraints on the models, and the predicted absolute fluxes, that are independent of the overall spectral shape. Further, the distances derived from scaling the model spectra to observed spectra on different dates and over different wavelength ranges provide a test of the self-consistency of the models. We find that the distances derived from the scale factors are well within the range of values for M82 ($D = 3.2\text{--}5.5$ Mpc with most values between $3.2\text{--}3.9$ Mpc; for example Tutui & Sofue 1997; Karachentsev & Kashibadze 2006; Dalcanton et al. 2009). The best fit reddened model spectra shown in Figure 3 yield distances of 3.7 Mpc (for the model corresponding to 33.15 days past explosion) and 4.0 Mpc (for the model corresponding to 44.12 days past explosion); the model spectra for 40.11 and 48.53 days past explosion yield distances of 4.2 and 3.5 Mpc,

respectively. The agreement of the values, both with each other and independent distance estimates, is fairly good, especially in light of the ($\sim 5\text{--}10\%$) uncertainties in the absolute flux levels of the observed spectra resulting from the scaling needed to match the interpolated photometric values. Finally, the comparison with the models also serves to illustrate the relative insensitivity of NIR observations to the effects of reddening, even for an object such as SN 2014J for which $A_V \sim 2.0$ mag (Amanullah et al. 2014; Foley et al. 2014).

The remarkably good agreement between the observed spectra and the DDC10 model spectra indicates that this model captures the basic elements of the physics involved in the formation and evolution of the observed NIR spectra of SNe Ia. The comparison between the model spectra and the observed spectra in the H and K band further supports the conclusions of Gall et al. (2012) and Friesen et al. (2014), who suggested that the post-maximum spectra of normal SNe Ia, including that of SN 2014J, longward of $\sim 1.5 \mu\text{m}$ are dominated by broad emission from permitted lines of Co II, but calls into question the previous identification of Co II absorption features at high velocity in the K -band spectra of other SNe Ia by Marion et al. (2009). It also suggests that the identification of the $1.77 \mu\text{m}$ feature as emission from Co II, dominated by the $1.7770 \mu\text{m}$ transition, is probably correct, despite the fact that the best fitting model spectrum predicts a somewhat wider and shifted line compared to the observations on February 19. The discrepancy between the observed and predicted line shape suggests that the relative strengths of the Co II lines contributing to the $1.77 \mu\text{m}$ feature are incorrect in the model. The $1.7463 \mu\text{m}$ line is far stronger in the model than the observations on this date indicate. As demonstrated in Figure 4, fitting the emission feature with multiple Gaussians, constrained to have the same velocity widths and centroids fixed at 1.746 , 1.777 , and $1.807 \mu\text{m}$, reveals that the observed line shape in the data from February 19 can be successfully

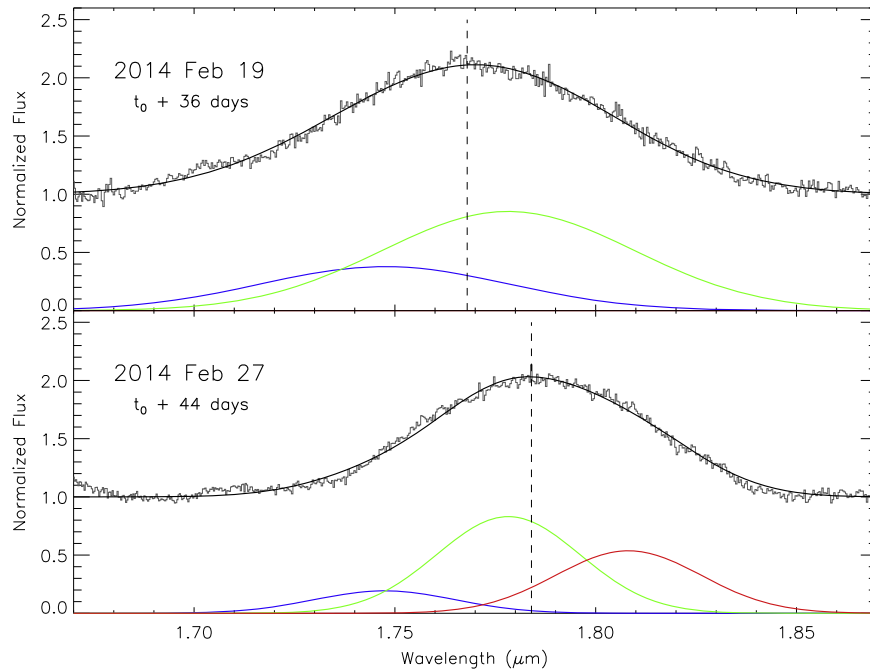


Figure 4. Decomposition of the observed $1.77\ \mu\text{m}$ feature in the spectra obtained on 2014 February 19 (top) and 2014 February 27 (bottom) into three Gaussians fixed at the locations of the strongest Co II lines ($1.7463\ \mu\text{m}$ —blue, $1.7770\ \mu\text{m}$ —green, and $1.8069\ \mu\text{m}$ —red). The solid black is the sum. The dashed lines mark the centroids of the observed feature. The observed shift in the feature is due to the changing line strengths of the three components.

reproduced with relative line intensities in the ratio of $\sim 0.45:1:0$. To reproduce the line shape observed in the February 27 data requires ratios of relative intensities of $\sim 0.25:1:0.65$. Therefore, the observed shift in the $1.77\ \mu\text{m}$ feature is due to the weakening of the $1.7463\ \mu\text{m}$ line and the strengthening of the $1.8069\ \mu\text{m}$ line. The velocity widths of the lines also decrease from $\sim 12,500\ \text{km s}^{-1}$ on February 19 to $\sim 7100\ \text{km s}^{-1}$ on February 27 as the ejecta expand and the optical depth decreases.

Following the suggestion of Gall et al. (2012), we used the width of the unblended Co II lines at 2.36 and $2.46\ \mu\text{m}$ to derive an estimate for the extent in velocity space of the Co-rich core. The weaker Co II line at $2.36\ \mu\text{m}$ seen in the February 27 data yields a velocity extent of $\sim 7000\ \text{km s}^{-1}$, in agreement with the line widths estimated above for the components of the $1.77\ \mu\text{m}$ feature. The stronger Co II $2.46\ \mu\text{m}$ line, the full width of which is generally inaccessible from the ground, yields a velocity extent of the Co core of $\sim 11,000\ \text{km s}^{-1}$ for both observation dates (February 19 and February 27), in good agreement with the location of the iron-group zone in models such as W7 and DDC10. Since almost all of the Co results from the decay of ^{56}Ni , the velocity extent of the Co should reflect the distribution of ^{56}Ni immediately after the explosion. As can be seen in Figure 1 of Dessart et al. (2014), the DDC10 model predicts that 99% of the ^{56}Ni should be located within $\sim 11,600\ \text{km s}^{-1}$ (Blondin et al. 2013, 2015). Our estimated velocity range is in good agreement with this value, as well as with the values found by Gall et al. (2012) for SN 2005cf and SN 2002bo at ~ 30 – 40 days past maximum B light (slightly later than our observations of SN 2014J at ~ 18 – 26 days past maximum B light).

Despite the success of the DDC10 model in reproducing the H - and K -band spectra of SN 2014J, the model does not match as well at other wavelengths. In Figure 5 we show the comparison between the FLITECAM spectra obtained on

February 25 and 27 and the DDC10 model over the entire 1 – $3.5\ \mu\text{m}$ range. Clear discrepancies between the data and the model are readily apparent in the J band, the region around $2\ \mu\text{m}$, and longwards of $\sim 3\ \mu\text{m}$. The disagreement between the model spectra and the observed spectrum in the $2\ \mu\text{m}$ region at 44 days past first light is somewhat surprising given how well the model matches the spectrum in this wavelength region earlier in the temporal evolution. Although Friesen et al. (2014) suggest that the $2\ \mu\text{m}$ feature seen in their spectra of a sample of SNe Ia (including SN 2014J) is due to emission from [Ni II] at $1.939\ \mu\text{m}$, the DDC10 model indicates that this feature is due primarily to [Co III] at $2.002\ \mu\text{m}$. In fact, very little ^{56}Ni should remain at this time, as 44 days past explosion is ~ 7 times longer than the 6 day half-life of ^{56}Ni , and therefore Friesen et al. (2014) speculated that the [Ni II] emission might arise from the stable isotope ^{58}Ni . However, the observed emission feature exhibits a peak at $\sim 1.99\ \mu\text{m}$, a wavelength which is far longer than that expected if [Ni II] were the primary component ($\sim 1.94\ \mu\text{m}$). Friesen et al. (2014) noted a similar wavelength discrepancy for this feature (relative to the predicted emission from [Ni II]) but attributed the wavelength difference to a redshift of the observed emission due to the combined effects of an asymmetrical explosion and Doppler shifts. Attributing the emission to [Co III] instead seems to be a far more satisfying explanation, especially given the success that spherical models have in reproducing the spectra. Nevertheless, the DDC10 model predicts this [Co III] feature decreases in strength far faster than the observations indicate; the emission peak is still very strong 44 days past first light, whereas the models predict the feature should have faded completely by about day 40.

In Figure 6 we present a comparison between the DDC10 model spectra and our J -band spectra. While the model seems to agree reasonably well with our observed J -band spectra at the earlier time in the evolution (~ 36 days past first light), the mismatch between the observations and the model at later times

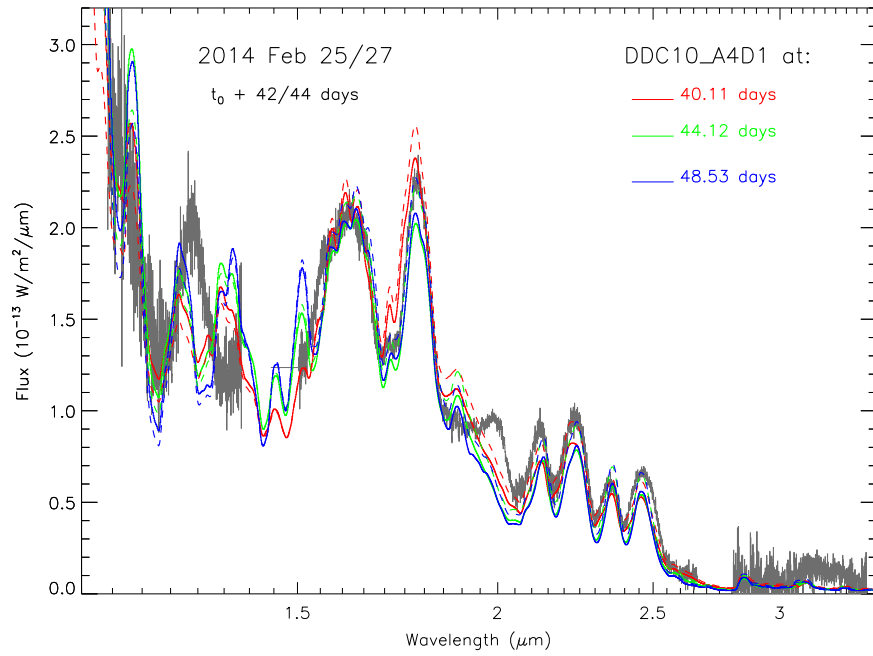


Figure 5. Comparison between the observed spectrum of SN 2014J obtained with FLITECAM on 2014 February 25 (2.8–3.4 μm) and February 27 (JHK bands) UT (in gray) and the CMFGEN model DDC10 A4D1 of Dessart et al. (2014) for three different times after explosion. The model spectra have been scaled to match the flux levels of the data. Dashed lines are model spectra that include reddening due to both Milky Way and M82 dust.

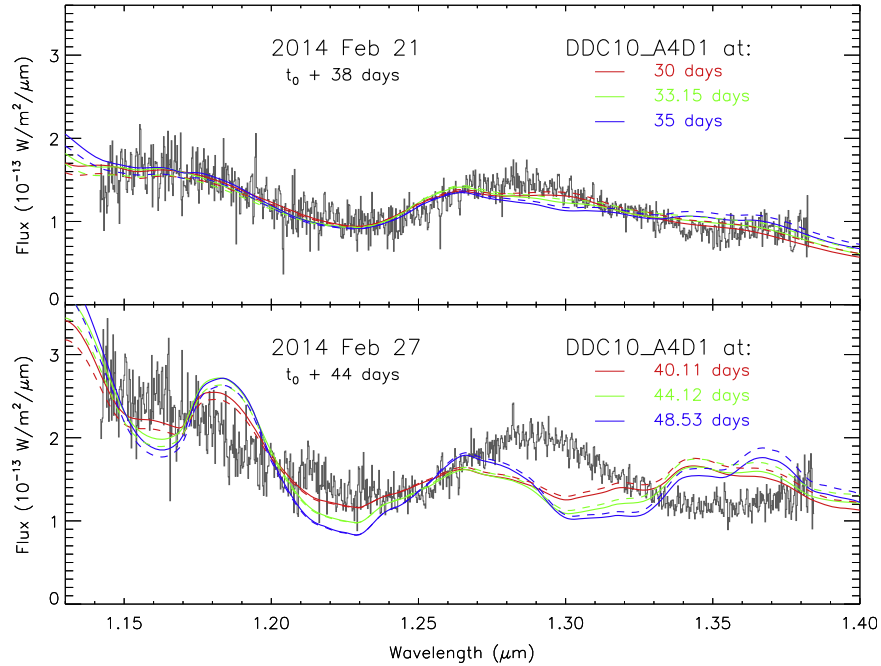


Figure 6. Comparison between the observed J -band spectrum of SN 2014J obtained with FLITECAM on 2014 February 21 UT (top) and 2014 February 27 UT (bottom) and the CMFGEN model DDC10 A4D1 of Dessart et al. (2014) for three different times past explosion. The model spectra have been scaled to match the flux levels of the data in this wavelength range. Dashed lines are model spectra that include reddening due to both Milky Way and M82 dust.

(~ 44 days past first light) is particularly bad. The most prominent characteristic of this wavelength region is a relatively strong emission peak at $\sim 1.29 \mu\text{m}$, present in both of our J -band spectra (Figure 2). The model spectra, particularly for late times in the temporal evolution of the SN, do not adequately reproduce this feature. Similar emission peaks can be seen in some of the late time spectra presented by Marion et al. (2009) and Gall et al. (2012), who attribute the emission to a blend of Fe II lines. Again, as for the $1.77 \mu\text{m}$ and

$2.46 \mu\text{m}$ features, the significantly higher atmospheric transmission afforded by *SOFIA* observations allows this $1.29 \mu\text{m}$ feature to be discerned and measured easily in the FLITECAM spectra.

The comparison between the observed 2.8–3.4 μm spectrum and the model spectra is shown in Figure 7, where a large discrepancy is also apparent, despite the low S/N of the data. (The low S/N of this spectrum is partially due to the additional background from the warm transfer optics in the dual

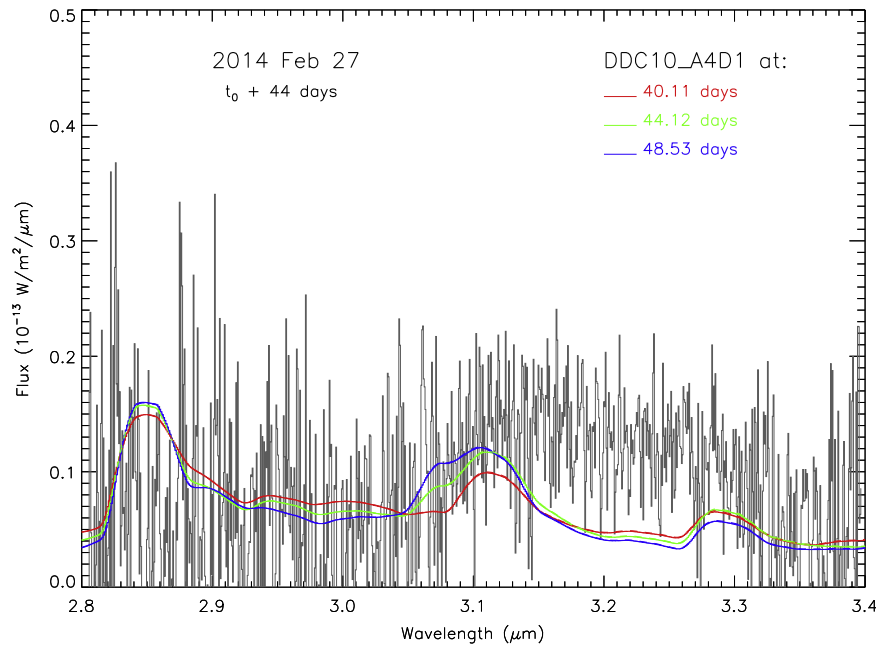


Figure 7. Comparison between the observed 2.8–3.4 μm spectrum of SN 2014J obtained with FLITECAM on 2014 February 27 UT and the CMFGEN model DDC10 A4D1 of Dessart et al. (2014) for three different times past explosion. The model spectra have been scaled to match the flux level of the data in this wavelength range.

instrument configuration.) The FLITECAM data indicate that broad emission is present between $\sim 3.0 \mu\text{m}$ and $3.4 \mu\text{m}$, which can provide additional constraints on models of SNe Ia. This spectral region is very difficult to observe from the ground due to strong telluric absorption, and our data may represent the first $\sim 3 \mu\text{m}$ spectra of a SN Ia ever obtained.

4. CONCLUSIONS

By providing access to wavelength regions that are difficult or impossible to observe from the ground, *SOFIA* can contribute significantly to studies of nearby bright SNe. We have demonstrated this with our FLITECAM observations of a normal SN Ia, SN 2014J in M82, about 36–44 days after first light, which have fully revealed a strong emission feature at $1.77 \mu\text{m}$. This feature shifted redward in wavelength and simultaneously narrowed between our two sets of observations separated by ~ 8 days. Because this feature extends into a wavelength regime that is difficult to observe from the ground, the identification, width, and the shift of the feature have never before been appreciated. Comparison with the non-LTE delayed detonation DDC10 A4D1 model of Dessart et al. (2014) suggests that this feature is most likely due to a blend of permitted emission lines from Co II. The variation in location and width is most likely due to changes in the relative strengths of the component lines. Other peaks seen in the *K*-band spectra are also due to permitted Co II lines emitting at the systemic velocity, as opposed to highly redshifted absorption features as suggested previously. The width of the Co II $2.46 \mu\text{m}$, a transition which is generally inaccessible from the ground, indicates that the Co-rich core extends to about $11,000 \text{ km s}^{-1}$ in agreement with predictions from models. The Dessart et al. (2014) model matches the observed spectra remarkably well across the 1.5 – $2.6 \mu\text{m}$ range, and scaling the model spectra to match the observed flux levels yields distances consistent with previous estimates for M82. However, the model provides a much poorer match to the observed spectra in the *J* band, and at

2 and $\sim 3 \mu\text{m}$ around 44 days after first light (~ 26 days after *B* maximum). The *J*-band spectra are characterized by a strong broad emission peak at $\sim 1.29 \mu\text{m}$, which is not reproduced in the model, particularly at late times in the temporal evolution. Similarly, the observed spectrum between 3.0 and $3.4 \mu\text{m}$, the first SN Ia spectrum ever obtained in this wavelength range, exhibits broad emission that is not seen in the model. The region around $2 \mu\text{m}$ exhibits a strong emission feature of [Co III] that decreases in strength far less quickly than the model predicts.

Our results serve to draw attention to the usefulness of the NIR regime, including the regions between the standard ground-based NIR passbands, for constraining SN models, due in part to the relative insensitivity of the NIR to the effects of reddening. The wavelength coverage afforded by FLITECAM on *SOFIA* allows the identification and study of the evolution of features not previously recognized in ground-based spectra of SNe Ia.

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